

Bootstrapping $\mathcal{N} \geq 2$ SCFTs

Madalena Lemos



String-Math 2017

July 26 2017

Based on:

1312.5344 w/ C. Beem, P. Liendo, W. Peelaers, L. Rastelli and B. van Rees

1511.07449 w/ P. Liendo

1702.05101 w/ M. Cornagliotto and V. Schomerus

Outline

① The (Super)conformal Bootstrap Program

Conformal bootstrap

Superconformal bootstrap

② A solvable subsector

③ Constraining the space of $\mathcal{N} = 2$ SCFTs

④ $4d$ $\mathcal{N} = 3$ SCFTs

⑤ Summary and Outlook

Outline

1 The (Super)conformal Bootstrap Program

Conformal bootstrap

Superconformal bootstrap

2 A solvable subsector

3 Constraining the space of $\mathcal{N} = 2$ SCFTs

4 $4d$ $\mathcal{N} = 3$ SCFTs

5 Summary and Outlook

The (Super)conformal Bootstrap Program

The (Super)conformal Bootstrap Program

What is the space of consistent (S)CFTs?

The (Super)conformal Bootstrap Program

What is the space of consistent (S)CFTs?

- Maximally supersymmetric theories: well known list of theories

The (Super)conformal Bootstrap Program

What is the space of consistent (S)CFTs?

- Maximally supersymmetric theories: well known list of theories

- $\mathcal{N} = 2$ theories: large known list of theories
many lacking a Lagrangian description

The (Super)conformal Bootstrap Program

What is the space of consistent (S)CFTs?

- Maximally supersymmetric theories: well known list of theories
- $\mathcal{N} = 3$ theories: not known to exist until García-Etxebarria and Regalado
- $\mathcal{N} = 2$ theories: large known list of theories many lacking a Lagrangian description

The (Super)conformal Bootstrap Program

What is the space of consistent (S)CFTs?

- Maximally supersymmetric theories: well known list of theories
- $\mathcal{N} = 3$ theories: not known to exist until García-Etxebarria and Regalado
- $\mathcal{N} = 2$ theories: large known list of theories many lacking a Lagrangian description

Can we bootstrap specific theories?

The (Super)conformal Bootstrap Program

What is the space of consistent (S)CFTs?

- Maximally supersymmetric theories: well known list of theories
- $\mathcal{N} = 3$ theories: not known to exist until García-Etxebarria and Regalado
- $\mathcal{N} = 2$ theories: large known list of theories many lacking a Lagrangian description

Can we bootstrap specific theories?

- Particularly helpful if theory is uniquely fixed by a set of discrete data

Outline

① The (Super)conformal Bootstrap Program

Conformal bootstrap

Superconformal bootstrap

② A solvable subsector

③ Constraining the space of $\mathcal{N} = 2$ SCFTs

④ $4d$ $\mathcal{N} = 3$ SCFTs

⑤ Summary and Outlook

Conformal Bootstrap

Conformal field theory defined by

Set of *local* operators and their correlation functions

Conformal Bootstrap

Conformal field theory defined by

Set of *local* operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

Conformal Bootstrap

Conformal field theory defined by

Set of *local* operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_k \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} \mathcal{O}_k(0)$$

Conformal Bootstrap

Conformal field theory defined by

Set of *local* operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_k \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} \mathcal{O}_k(0)$$

→ Finite radius of convergence

Conformal Bootstrap

Conformal field theory defined by

Set of *local* operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k\text{prim.}} \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} c(x, \partial_x)\mathcal{O}_k(0)$$

→ Finite radius of convergence

Conformal Bootstrap

Conformal field theory defined by

Set of *local* operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k\text{prim.}} \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} c(x, \partial_x)\mathcal{O}_k(0)$$

→ Finite radius of convergence

→ n -point function by recursive use of the OPE until

$$\langle \mathbb{1} \rangle = 1$$

Conformal Bootstrap

Conformal field theory defined by

Set of *local* operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\{\lambda_{\mathcal{O}_i\mathcal{O}_j\mathcal{O}_k}\}$

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k\text{prim.}} \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} c(x, \partial_x) \mathcal{O}_k(0)$$

→ Finite radius of convergence

→ n -point function by recursive use of the OPE until

$$\langle \mathbb{1} \rangle = 1$$

Conformal Bootstrap

Conformal field theory defined by

Set of *local* operators and their correlation functions

CFT data

$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\{\lambda_{\mathcal{O}_i\mathcal{O}_j\mathcal{O}_k}\}$

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k\text{prim.}} \lambda_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} c(x, \partial_x)\mathcal{O}_k(0)$$

- Finite radius of convergence
- n -point function by recursive use of the OPE until $\langle \mathbb{1} \rangle = 1$

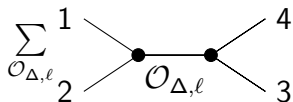
CFT data strongly constrained

- ▶ Unitarity
- ▶ Associativity of the operator product algebra

Conformal Bootstrap

Crossing Symmetry

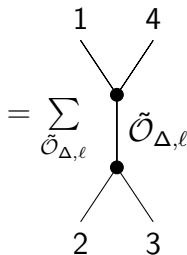
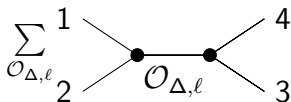
$$\langle (\mathcal{O}_1(x_1) \mathcal{O}_2(x_2)) \mathcal{O}_3(x_3) \mathcal{O}_4(x_4) \rangle =$$



Conformal Bootstrap

Crossing Symmetry

$$\langle \mathcal{O}_1(x_1) (\mathcal{O}_2(x_2) \mathcal{O}_3(x_3)) \mathcal{O}_4(x_4) \rangle =$$



Outline

1 The (Super)conformal Bootstrap Program

Conformal bootstrap

Superconformal bootstrap

2 A solvable subsector

3 Constraining the space of $\mathcal{N} = 2$ SCFTs

4 $4d$ $\mathcal{N} = 3$ SCFTs

5 Summary and Outlook

The Superconformal Bootstrap

- ▶ Various conformal families related by action of supercharges

The Superconformal Bootstrap

- ▶ Various conformal families related by action of supercharges
- ▶ Finite re-organization of an infinite amount of data

The Superconformal Bootstrap

- ▶ Various conformal families related by action of supercharges
- ▶ Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

The Superconformal Bootstrap

- ▶ Various conformal families related by action of supercharges
- ▶ Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

→ Yes, for $4d \mathcal{N} \geq 2$ [Beem ML Liendo Peelaers Rastelli van Rees]

The Superconformal Bootstrap

- ▶ Various conformal families related by action of supercharges
- ▶ Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

→ Yes, for $4d \mathcal{N} \geq 2$ [Beem ML Liendo Peelaers Rastelli van Rees]
 $6d \mathcal{N} = (2, 0)$ and $2d \mathcal{N} = (0, 4)$ [Beem Rastelli van Rees]

- ▶ Step 1: Solve this subsector
- ▶ (Step 2: Full blown numerics for the rest)

Outline

① The (Super)conformal Bootstrap Program

Conformal bootstrap

Superconformal bootstrap

② A solvable subsector

③ Constraining the space of $\mathcal{N} = 2$ SCFTs

④ $4d$ $\mathcal{N} = 3$ SCFTs

⑤ Summary and Outlook

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{ \mathcal{O}_{\Delta, (j_1, j_2)}, \quad \}$$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\left\{ \mathcal{O}_{\Delta, (j_1, j_2)}, \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f \right\}$$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\left\{ \mathcal{O}_{\Delta, (j_1, j_2)}, \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f \right\}$$

Claim

→ Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\left\{ \mathcal{O}_{\Delta, (j_1, j_2)}, \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f \right\}$$

Claim

- Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$
- Restrict to operators with $\Delta = 2R + j_1 + j_2$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\left\{ \mathcal{O}_{\Delta, (j_1, j_2)}, \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f \right\}$$

Claim

- Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$
- Restrict to operators with $\Delta = 2R + j_1 + j_2$

$$\langle \mathcal{O}_1^{I_1}(z_1, \bar{z}_1) \dots \mathcal{O}_n^{I_n}(z_n, \bar{z}_n) \rangle$$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1, j_2)}, \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f\}$$

Claim

- Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$
- Restrict to operators with $\Delta = 2R + j_1 + j_2$

$$u_{l_1}(\bar{z}_1) \dots u_{l_n}(\bar{z}_n) \langle \mathcal{O}_1^{l_1}(z_1, \bar{z}_1) \dots \mathcal{O}_n^{l_n}(z_n, \bar{z}_n) \rangle$$

Chiral algebra

Organize operators in representations of superconformal algebra

$$\{\mathcal{O}_{\Delta, (j_1, j_2)}, \underbrace{R}_{SU(2)_R}, \underbrace{r}_{U(1)_r}, f\}$$

Claim

→ Pick a plane $\mathbb{R}^2 \in \mathbb{R}^4$, $(z, \bar{z}) \in \mathbb{R}^2$

→ Restrict to operators with $\Delta = 2R + j_1 + j_2$

$$u_{l_1}(\bar{z}_1) \dots u_{l_n}(\bar{z}_n) \langle \mathcal{O}_1^{l_1}(z_1, \bar{z}_1) \dots \mathcal{O}_n^{l_n}(z_n, \bar{z}_n) \rangle = f(z_i)$$

→ Meromorphic!

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent \mathbb{Q}

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- Cohomology at the origin \Rightarrow non-empty classes

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
 - Cohomology at the origin \Rightarrow non-empty classes
- $$\Delta = 2R + j_1 + j_2$$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $Q \sim Q + S$

→ Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

- ▶ On plane $\underbrace{\mathfrak{sl}_2} \times \underbrace{\bar{\mathfrak{sl}}_2}$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$

→ Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

- ▶ On plane $\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}$

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$

→ Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

- ▶ On plane $\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}_{\text{does not}}$

Chiral algebra

Why?

▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$

→ Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

▶ On plane $\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}_{\text{does not}}$

→ twisted translations $u_I(\bar{z})$

Chiral algebra

Why?

▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$

→ Cohomology at the origin \Rightarrow non-empty classes

$$\Delta = 2R + j_1 + j_2$$

▶ On plane $\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}_{\text{does not}}$

→ twisted translations $u_I(\bar{z})$

↪ diagonal subalgebra $\bar{\mathfrak{sl}}_2 \times \mathfrak{su}(2)_R$ is \mathbb{Q} exact

Chiral algebra

Why?

- ▶ Subsector = Cohomology of nilpotent $\mathbb{Q} \sim \mathcal{Q} + \mathcal{S}$
- Cohomology at the origin \Rightarrow non-empty classes
 $\Delta = 2R + j_1 + j_2$
- ▶ On plane $\underbrace{\mathfrak{sl}_2}_{\text{commutes with } \mathbb{Q}} \times \underbrace{\bar{\mathfrak{sl}}_2}_{\text{does not}}$
- twisted translations $u_I(\bar{z})$
- ↪ diagonal subalgebra $\bar{\mathfrak{sl}}_2 \times \mathfrak{su}(2)_R$ is \mathbb{Q} exact
- ↪ anti-holomorphic dependence drops out

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

\rightarrow Cohomology classes \Rightarrow Vertex operators

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

- \rightarrow Cohomology classes \Rightarrow Vertex operators
- \rightarrow conformal weight $h = R + j_1 + j_2 \geq 0$

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

- \rightarrow Cohomology classes \Rightarrow Vertex operators
- \rightarrow conformal weight $h = R + j_1 + j_2 \geq 0$
- \rightarrow Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

- \rightarrow Cohomology classes \Rightarrow Vertex operators
- \rightarrow conformal weight $h = R + j_1 + j_2 \geq 0$
- \rightarrow Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary
 - \blacktriangleright Very specific non-unitary VOA constrained by unitarity of $4d$ theory

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

- \rightarrow Cohomology classes \Rightarrow Vertex operators
- \rightarrow conformal weight $h = R + j_1 + j_2 \geq 0$
- \rightarrow Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary
 - \blacktriangleright Very specific non-unitary VOA constrained by unitarity of $4d$ theory
 - \rightarrow some operators acquire negative norms

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

- \rightarrow Cohomology classes \Rightarrow Vertex operators
- \rightarrow conformal weight $h = R + j_1 + j_2 \geq 0$
- \rightarrow Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary
 - \blacktriangleright Very specific non-unitary VOA constrained by unitarity of $4d$ theory
 - \rightarrow some operators acquire negative norms
 - \blacktriangleright Obtain VOA from $4d \mathcal{N} = 2$ SCFT

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

- \rightarrow Cohomology classes \Rightarrow Vertex operators
- \rightarrow conformal weight $h = R + j_1 + j_2 \geq 0$
- \rightarrow Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary
 - \blacktriangleright Very specific non-unitary VOA constrained by unitarity of $4d$ theory
 - \rightarrow some operators acquire negative norms
 - \blacktriangleright Obtain VOA from $4d \mathcal{N} = 2$ SCFT
 - \blacktriangleright Given a VOA does there exist a $4d$ SCFT?

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

- \rightarrow Cohomology classes \Rightarrow Vertex operators
- \rightarrow conformal weight $h = R + j_1 + j_2 \geq 0$
- \rightarrow Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary
 - \blacktriangleright Very specific non-unitary VOA constrained by unitarity of $4d$ theory
 - \rightarrow some operators acquire negative norms
 - \blacktriangleright Obtain VOA from $4d \mathcal{N} = 2$ SCFT
 - \blacktriangleright Given a VOA does there exist a $4d$ SCFT?
 - \hookrightarrow give an example of what can go wrong

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

- \rightarrow Cohomology classes \Rightarrow Vertex operators
- \rightarrow conformal weight $h = R + j_1 + j_2 \geq 0$
- \rightarrow Each $\mathcal{N} = 2$ multiplet contributes at most with one \mathfrak{sl}_2 primary
 - \blacktriangleright Very specific non-unitary VOA constrained by unitarity of $4d$ theory
 - \rightarrow some operators acquire negative norms
 - \blacktriangleright Obtain VOA from $4d \mathcal{N} = 2$ SCFT
 - \blacktriangleright Given a VOA does there exist a $4d$ SCFT?
 - \hookrightarrow give an example of what can go wrong
 - \blacktriangleright How much information can we recover from the VOA?

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

\rightarrow Stress tensor $T_{\mu\nu}$

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

\rightarrow Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

- \rightarrow Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- \rightarrow Stress tensor supermultiplet

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

- \rightarrow Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- \rightarrow Stress tensor supermultiplet $\Rightarrow 2d$ stress tensor

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

- \rightarrow Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- \rightarrow Stress tensor supermultiplet $\Rightarrow 2d$ stress tensor

$$T(z)T(0) \sim -12\frac{c_{4d}/2}{z^4} + 2\frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

- \rightarrow Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- \rightarrow Stress tensor supermultiplet $\Rightarrow 2d$ stress tensor

$$T(z)T(0) \sim -12\frac{c_{4d}/2}{z^4} + 2\frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

- \hookrightarrow Global \mathfrak{sl}_2 enhances to Virasoro

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

- \rightarrow Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- \rightarrow Stress tensor supermultiplet $\Rightarrow 2d$ stress tensor

$$T(z)T(0) \sim -12\frac{c_{4d}/2}{z^4} + 2\frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

\hookrightarrow Global \mathfrak{sl}_2 enhances to Virasoro

\hookrightarrow $c_{2d} = -12c_{4d}$

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

- \rightarrow Stress tensor $T_{\mu\nu} \rightsquigarrow$ superdescendant
- \rightarrow Stress tensor supermultiplet $\Rightarrow 2d$ stress tensor

$$T(z)T(0) \sim -12\frac{c_{4d}/2}{z^4} + 2\frac{T(0)}{z^2} + \frac{\partial T(0)}{z} + \dots,$$

- \hookrightarrow Global \mathfrak{sl}_2 enhances to Virasoro
- \hookrightarrow $c_{2d} = -12c_{4d}$
- \hookrightarrow Virasoro representations seem to mix different types of $4d$ multiplets

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

\rightarrow Flavor symmetries current multiplet

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

- \rightarrow Flavor symmetries current multiplet
- \hookrightarrow Affine Kac Moody current algebra

$$J^a(z)J^b(0) \sim -\frac{k_{4d}/2\delta^{ab}}{z^2} + if^{abc}\frac{J^c(0)}{z} + \dots,$$

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

- \rightarrow Flavor symmetries current multiplet
- \hookrightarrow Affine Kac Moody current algebra

$$J^a(z)J^b(0) \sim -\frac{k_{4d}/2\delta^{ab}}{z^2} + if^{abc}\frac{J^c(0)}{z} + \dots,$$

$$\hookrightarrow k_{2d} = -\frac{k_{4d}}{2}$$

$4d \mathcal{N} \geq 2$ SCFT \longrightarrow VOA

Which operators are in the cohomology?

- \rightarrow Flavor symmetries current multiplet
- \hookrightarrow Affine Kac Moody current algebra

$$J^a(z)J^b(0) \sim -\frac{k_{4d}/2\delta^{ab}}{z^2} + if^{abc}\frac{J^c(0)}{z} + \dots,$$

\hookrightarrow

$$k_{2d} = -\frac{k_{4d}}{2}$$

$\rightarrow \dots$

What is the space of consistent SCFTs?

4d $\mathcal{N} = 2$ SCFTs with a flavor symmetry

$\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$

What is the space of consistent SCFTs?

$4d \mathcal{N} = 2$ SCFTs with a flavor symmetry

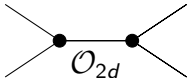
$\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$ functions of c_{2d} and k_{2d}

What is the space of consistent SCFTs?

$4d \mathcal{N} = 2$ SCFTs with a flavor symmetry

$\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$ functions of c_{2d} and k_{2d}

- ▶ Block decomposition:

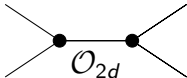
$$\sum_{\mathcal{O}_{2d}} \lambda_{\mathcal{O}_{2d}}^2 \text{ (diagram) }$$


What is the space of consistent SCFTs?

$4d \mathcal{N} = 2$ SCFTs with a flavor symmetry

$\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$ functions of c_{2d} and k_{2d}

- ▶ Block decomposition:

$$\sum_{\mathcal{O}_{2d}} \lambda_{\mathcal{O}_{2d}}^2 \text{ (diagram) }$$


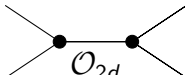
$$\rightarrow \lambda_{\mathcal{O}_{2d}}^2$$

What is the space of consistent SCFTs?

$4d \mathcal{N} = 2$ SCFTs with a flavor symmetry

$\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$ functions of c_{2d} and k_{2d}

- ▶ Block decomposition:

$$\sum_{\mathcal{O}_{2d}} \lambda_{\mathcal{O}_{2d}}^2 \text{ (diagram) }$$


$$\rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2$$

Interpret as four-dimensional quantities

(with some assumptions: interacting theory, unique stress tensor)

What is the space of consistent SCFTs?

$4d \mathcal{N} = 2$ SCFTs with a flavor symmetry

$\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$ functions of c_{2d} and k_{2d}

- ▶ Block decomposition:

$$\sum_{\mathcal{O}_{2d}} \lambda_{\mathcal{O}_{2d}}^2 \text{ [diagram of a block with two internal vertices and four external legs]} \mathcal{O}_{2d}$$

$$\rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2 \underbrace{\geq}_{4d \text{ unitarity}} 0$$

Interpret as four-dimensional quantities

(with some assumptions: interacting theory, unique stress tensor)

What is the space of consistent SCFTs?

$4d \mathcal{N} = 2$ SCFTs with a flavor symmetry

$\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$ functions of c_{2d} and k_{2d}

- ▶ Block decomposition:

$$\sum_{\mathcal{O}_{2d}} \lambda_{\mathcal{O}_{2d}}^2 \text{ (diagram) }$$

The diagram shows a central horizontal line segment with two black dots at its ends. From each dot, two lines extend outwards at an angle, forming a total of four external lines. Below the central line segment is the label \mathcal{O}_{2d} .

$$\rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2 \underbrace{\geq}_{4d \text{ unitarity}} 0 \Rightarrow \text{New unitarity bounds}$$

Interpret as four-dimensional quantities

(with some assumptions: interacting theory, unique stress tensor)

Outline

① The (Super)conformal Bootstrap Program

Conformal bootstrap

Superconformal bootstrap

② A solvable subsector

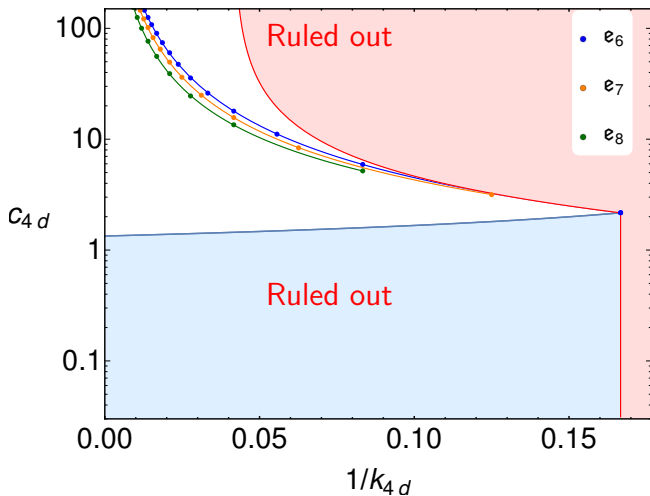
③ Constraining the space of $\mathcal{N} = 2$ SCFTs

④ $4d$ $\mathcal{N} = 3$ SCFTs

⑤ Summary and Outlook

$4d \mathcal{N} = 2$ SCFTs with E_6 flavor symmetry

$$\langle TTTT \rangle, \langle J^a J^b J^c J^d \rangle, \langle TTJ^a J^b \rangle$$



Outline

① The (Super)conformal Bootstrap Program

Conformal bootstrap

Superconformal bootstrap

② A solvable subsector

③ Constraining the space of $\mathcal{N} = 2$ SCFTs

④ $4d$ $\mathcal{N} = 3$ SCFTs

⑤ Summary and Outlook

$\mathcal{N} = 3$ Chiral algebra

- ▶ $4d \mathcal{N} \geq 3$: some of the extra supercharges commute with \mathbb{Q}

$\mathcal{N} = 3$ Chiral algebra

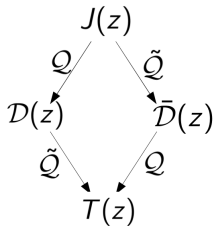
- ▶ $4d \mathcal{N} \geq 3$: some of the extra supercharges commute with \mathbb{Q}
 - ↳ $4d \mathcal{N} = 4 \Rightarrow 2d$ “small” $\mathcal{N} = 4$ chiral algebra

$\mathcal{N} = 3$ Chiral algebra

- ▶ $4d \mathcal{N} \geq 3$: some of the extra supercharges commute with \mathbb{Q}
 - ↪ $4d \mathcal{N} = 4 \Rightarrow 2d$ “small” $\mathcal{N} = 4$ chiral algebra
 - ↪ $4d \mathcal{N} = 3 \Rightarrow 2d \mathcal{N} = 2$ chiral algebra [Nishinaka, Tachikawa]

$\mathcal{N} = 3$ Chiral algebra

- ▶ $4d \mathcal{N} \geq 3$: some of the extra supercharges commute with \mathbb{Q}
 - ↪ $4d \mathcal{N} = 4 \Rightarrow 2d$ “small” $\mathcal{N} = 4$ chiral algebra
 - ↪ $4d \mathcal{N} = 3 \Rightarrow 2d \mathcal{N} = 2$ chiral algebra [Nishinaka, Tachikawa]
- ▶ $2d$ stress tensor promoted to supermultiplet \mathcal{J}

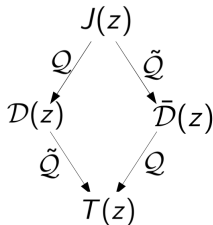


$\mathcal{N} = 3$ Chiral algebra

- ▶ $4d \mathcal{N} \geq 3$: some of the extra supercharges commute with \mathbb{Q}
 - ↪ $4d \mathcal{N} = 4 \Rightarrow 2d$ “small” $\mathcal{N} = 4$ chiral algebra
 - ↪ $4d \mathcal{N} = 3 \Rightarrow 2d \mathcal{N} = 2$ chiral algebra [Nishinaka, Tachikawa]
- ▶ $2d$ stress tensor promoted to supermultiplet \mathcal{J}

$2d \mathcal{N} = 2$ **Stress tensor** \mathcal{J}

→ A trivial statement in $2d$:



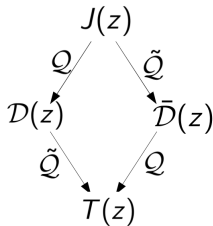
$\mathcal{N} = 3$ Chiral algebra

- ▶ $4d \mathcal{N} \geq 3$: some of the extra supercharges commute with \mathbb{Q}
 - ↪ $4d \mathcal{N} = 4 \Rightarrow 2d$ “small” $\mathcal{N} = 4$ chiral algebra
 - ↪ $4d \mathcal{N} = 3 \Rightarrow 2d \mathcal{N} = 2$ chiral algebra [Nishinaka, Tachikawa]
- ▶ $2d$ stress tensor promoted to supermultiplet \mathcal{J}

$2d \mathcal{N} = 2$ **Stress tensor** \mathcal{J}

→ A trivial statement in $2d$:

$\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J} \rangle$ is fixed in terms of c_{2d}



$\mathcal{N} = 3$ Chiral algebra

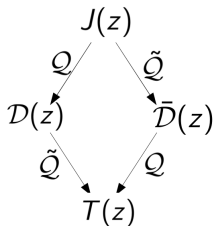
- ▶ $4d \mathcal{N} \geq 3$: some of the extra supercharges commute with \mathbb{Q}
 - ↪ $4d \mathcal{N} = 4 \Rightarrow 2d$ “small” $\mathcal{N} = 4$ chiral algebra
 - ↪ $4d \mathcal{N} = 3 \Rightarrow 2d \mathcal{N} = 2$ chiral algebra [Nishinaka, Tachikawa]
- ▶ $2d$ stress tensor promoted to supermultiplet \mathcal{J}

$2d \mathcal{N} = 2$ Stress tensor \mathcal{J}

→ A trivial statement in $2d$:

$\langle \mathcal{J} \mathcal{J} \mathcal{J} \mathcal{J} \rangle$ is fixed in terms of c_{2d}

→ Present in any local $\mathcal{N} = 3$ SCFT



Space of $\mathcal{N} = 3$ SCFTs

$2d$ $\mathcal{N} = 2$ **Stress tensor** \mathcal{J}

$\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J} \rangle$ is fixed in terms of $c_{2d} \Rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2$

Space of $\mathcal{N} = 3$ SCFTs

$2d$ $\mathcal{N} = 2$ **Stress tensor** \mathcal{J}

$\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J} \rangle$ is fixed in terms of $c_{2d} \Rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2$

$$c_{4d} \geq \frac{13}{24}$$

[Cornagliotto, ML, Schomerus]

\hookrightarrow Not saturated by any known SCFT

Space of $\mathcal{N} = 3$ SCFTs

$2d$ $\mathcal{N} = 2$ **Stress tensor** \mathcal{J}

$\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J} \rangle$ is fixed in terms of $c_{2d} \Rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2$

$$c_{4d} \geq \frac{13}{24}$$

[Cornagliotto, ML, Schomerus]

\hookrightarrow Not saturated by any known SCFT

Space of $\mathcal{N} = 3$ SCFTs

$2d$ $\mathcal{N} = 2$ **Stress tensor** \mathcal{J}

$\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J} \rangle$ is fixed in terms of $c_{2d} \Rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2$

$$c_{4d} \geq \frac{13}{24}$$

[Cornagliotto, ML, Schomerus]

- ↪ Not saturated by any known SCFT
- ↪ Similar bounds in $\mathcal{N} = 4$ and $\mathcal{N} = 2$ saturated by known SCFTs [Beem, Rastelli, van Rees] [Liendo, Ramirez, Seo]

Space of $\mathcal{N} = 3$ SCFTs

$2d$ $\mathcal{N} = 2$ **Stress tensor** \mathcal{J}

$\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J} \rangle$ is fixed in terms of $c_{2d} \Rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2$

$$c_{4d} \geq \frac{13}{24}$$

[Cornagliotto, ML, Schomerus]

- ↪ Not saturated by any known SCFT
- ↪ Similar bounds in $\mathcal{N} = 4$ and $\mathcal{N} = 2$ saturated by known SCFTs [Beem, Rastelli, van Rees] [Liendo, Ramirez, Seo]
- $c_{4d} = \frac{13}{24} \Rightarrow$ reconstruct $4d$ operators appearing in $\mathcal{J}\mathcal{J}$

Space of $\mathcal{N} = 3$ SCFTs

$2d$ $\mathcal{N} = 2$ **Stress tensor** \mathcal{J}

$\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J} \rangle$ is fixed in terms of $c_{2d} \Rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2$

$$c_{4d} \geq \frac{13}{24}$$

[Cornagliotto, ML, Schomerus]

- ↪ Not saturated by any known SCFT
- ↪ Similar bounds in $\mathcal{N} = 4$ and $\mathcal{N} = 2$ saturated by known SCFTs [Beem, Rastelli, van Rees] [Liendo, Ramirez, Seo]
- $c_{4d} = \frac{13}{24} \Rightarrow$ reconstruct $4d$ operators appearing in $\mathcal{J}\mathcal{J}$
- Signs of norms inconsistent with an *interacting* $4d$ SCFT existing

Space of $\mathcal{N} = 3$ SCFTs

$2d$ $\mathcal{N} = 2$ **Stress tensor** \mathcal{J}

$\langle \mathcal{J}\mathcal{J}\mathcal{J}\mathcal{J} \rangle$ is fixed in terms of $c_{2d} \Rightarrow \lambda_{\mathcal{O}_{2d}}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2$

$$c_{4d} > \frac{13}{24}$$

[Cornagliotto, ML, Schomerus]

- ↪ Not saturated by any known SCFT
- ↪ Similar bounds in $\mathcal{N} = 4$ and $\mathcal{N} = 2$ saturated by known SCFTs [Beem, Rastelli, van Rees] [Liendo, Ramirez, Seo]
- $c_{4d} = \frac{13}{24} \Rightarrow$ reconstruct $4d$ operators appearing in $\mathcal{J}\mathcal{J}$
- Signs of norms inconsistent with an *interacting* $4d$ SCFT existing

Outline

① The (Super)conformal Bootstrap Program

Conformal bootstrap

Superconformal bootstrap

② A solvable subsector

③ Constraining the space of $\mathcal{N} = 2$ SCFTs

④ $4d$ $\mathcal{N} = 3$ SCFTs

⑤ Summary and Outlook

Summary and Outlook

New constraints on the space of allowed $\mathcal{N} = 2, 3$ SCFTs

Summary and Outlook

New constraints on the space of allowed $\mathcal{N} = 2, 3$ SCFTs

→ No “minimal” $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$

Summary and Outlook

New constraints on the space of allowed $\mathcal{N} = 2, 3$ SCFTs

- No “minimal” $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$
 - ↪ can we improve on this bound analytically?
What are the conditions for a VOA to correspond to a $4d$ SCFT?

Summary and Outlook

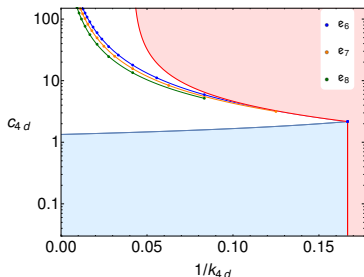
New constraints on the space of allowed $\mathcal{N} = 2, 3$ SCFTs

- No “minimal” $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$
 - ↪ can we improve on this bound analytically?
What are the conditions for a VOA to correspond to a $4d$ SCFT?
- Can the numerical bootstrap complement these?

Summary and Outlook

New constraints on the space of allowed $\mathcal{N} = 2, 3$ SCFTs

- No “minimal” $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$
 - ↳ can we improve on this bound analytically?
What are the conditions for a VOA to correspond to a $4d$ SCFT?
- Can the numerical bootstrap complement these?
- Is $c_{4d}/k_{4d} \geq \dots$?



Summary and Outlook

New constraint on the space of allowed $\mathcal{N} = 2, 3$ SCFTs

- No “minimal” $\mathcal{N} = 3$ SCFT with $c = \frac{13}{24}$
 - ↪ can we improve on this bound analytically?
What are the conditions for a VOA to correspond to a $4d$ SCFT?
- Can the numerical bootstrap complement these?
- Is $c_{4d}/k_{4d} \geq \dots$?

Numerically solving theories?

- ▶ This mixed correlator seems like a good starting point

Thank you!

Backup slides

Outline

Numerical conformal Bootstrap review

Chiral algebra

$4d \mathcal{N} = 3$ SCFTs

Solving $\mathcal{N} = 3$ SCFTs?

Conformal Bootstrap

→ Solve crossing equations for *all* four-point functions

Conformal Bootstrap

→ Solve crossing equations for *all* four-point functions

[Rattazzi Rychkov Tonni Vichi]

- ▶ Solving \Rightarrow constraining

Conformal Bootstrap

→ Solve crossing equations for *all* four-point functions

[Rattazzi Rychkov Tonni Vichi]

▶ Solving \Rightarrow constraining

→ Guess for the spectrum

Conformal Bootstrap

→ Solve crossing equations for *all* four-point functions

[Rattazzi Rychkov Tonni Vichi]

▶ Solving \Rightarrow constraining

→ Guess for the spectrum

↪ there's a large gap in the spectrum

Conformal Bootstrap

→ Solve crossing equations for *all* four-point functions

[Rattazzi Rychkov Tonni Vichi]

- ▶ Solving \Rightarrow constraining
 - Guess for the spectrum
 - \hookrightarrow there's a large gap in the spectrum
 - Can it ever define a consistent CFT?

Conformal Bootstrap

→ Solve crossing equations for *all* four-point functions

[Rattazzi Rychkov Tonni Vichi]

▶ Solving \Rightarrow constraining

→ Guess for the spectrum

↪ there's a large gap in the spectrum

→ Can it ever define a consistent CFT?

Sum rule: identical scalars ϕ

Conformal Bootstrap

→ Solve crossing equations for *all* four-point functions

[Rattazzi Rychkov Tonni Vichi]

- ▶ Solving \Rightarrow constraining
 - Guess for the spectrum
 - \hookrightarrow there's a large gap in the spectrum
 - Can it ever define a consistent CFT?

Sum rule: identical scalars ϕ

→ Identity operator $\lambda_{\mathcal{O}\mathcal{O}\mathbb{1}} = 1$

$$1 = \sum_{\substack{\mathcal{O}_{\Delta,\ell} \neq \mathbb{1} \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 \underbrace{\frac{u^{\Delta_\phi} g_{\Delta,\ell}(v, u) - v^{\Delta_\phi} g_{\Delta,\ell}(u, v)}{v^{\Delta_\phi} - u^{\Delta_\phi}}}_{F_{\Delta,\ell}}$$

Conformal Bootstrap

Sum rule

$$1 = \sum_{\substack{\mathcal{O}_{\Delta,\ell} \neq 1 \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}$$

Conformal Bootstrap

Sum rule

$$1 = \sum_{\substack{\mathcal{O}_{\Delta,\ell} \neq 1 \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}$$

- ▶ Find Functional Ψ such that
 - ↪ $\psi \cdot 1 < 0$ ($\mathbb{1}$)
 - ↪ $\psi \cdot F_{\Delta,\ell}(u, v) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum

Conformal Bootstrap

Sum rule

$$1 = \sum_{\substack{\mathcal{O}_{\Delta,\ell} \neq 1 \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}$$

► Find Functional Ψ such that

↪ $\psi \cdot \mathbb{1} < 0$ ($\mathbb{1}$)

↪ $\psi \cdot F_{\Delta,\ell}(u, v) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum

→ Spectrum is inconsistent \Rightarrow rule out CFT

Conformal Bootstrap

Sum rule

$$1 = \sum_{\substack{\mathcal{O}_{\Delta,\ell} \neq 1 \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}$$

- ▶ Find Functional Ψ such that

- ↪ $\psi \cdot 1 < 0$ ($\mathbb{1}$)

- ↪ $\psi \cdot F_{\Delta,\ell}(u, v) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum

→ Spectrum is inconsistent \Rightarrow rule out CFT

- ▶ Truncate

$$\psi = \sum_{m,n}^{m,n \leq \Lambda} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{z=\bar{z}=\frac{1}{2}}$$

Conformal Bootstrap

Sum rule

$$1 = \sum_{\substack{\mathcal{O}_{\Delta,\ell} \neq \mathbb{1} \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}$$

- ▶ Find Functional Ψ such that

- ↪ $\psi \cdot \mathbb{1} < 0$ ($\mathbb{1}$)

- ↪ $\psi \cdot F_{\Delta,\ell}(u, v) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum

→ Spectrum is inconsistent \Rightarrow rule out CFT

- ▶ Truncate

$$\psi = \sum_{\substack{m,n \leq \Lambda \\ m,n}} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{z=\bar{z}=\frac{1}{2}}$$

→ Increase $\Lambda \Rightarrow$ bounds get stronger

Conformal Bootstrap

Sum rule

$$1 = \sum_{\substack{\mathcal{O}_{\Delta,\ell} \neq 1 \\ \mathcal{O} \in \phi\phi}} \lambda_{\phi\phi\mathcal{O}}^2 F_{\Delta,\ell}$$

- ▶ Find Functional Ψ such that

- ↪ $\psi \cdot 1 < 0$ ($\mathbb{1}$)

- ↪ $\psi \cdot F_{\Delta,\ell}(u, v) \geq 0$ for all $\{\Delta, \ell\}$ in spectrum

→ Spectrum is inconsistent \Rightarrow rule out CFT

- ▶ Truncate

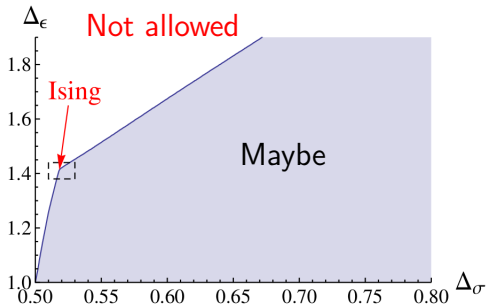
$$\psi = \sum_{\substack{m,n \leq \Lambda \\ m,n}} a_{mn} \partial_z^m \partial_{\bar{z}}^n |_{z=\bar{z}=\frac{1}{2}}$$

→ Increase $\Lambda \Rightarrow$ bounds get stronger

→ Always true bounds

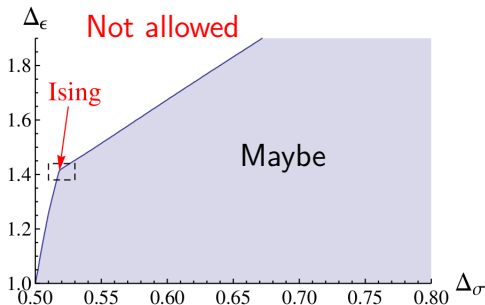
3d Ising Model

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, PRD 86 025022]



3d Ising Model

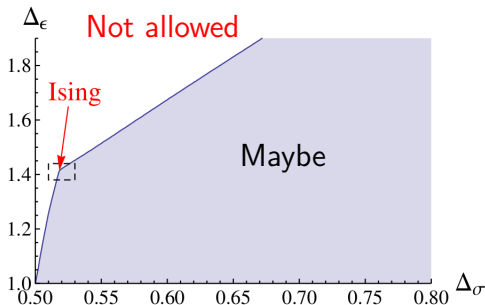
[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, PRD 86 025022]



→ Saturated by 3d Ising model

3d Ising Model

[El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin, Vichi, PRD 86 025022]



- Saturated by 3d Ising model
- 3d Ising lives at “kink”

Outline

Numerical conformal Bootstrap review

Chiral algebra

$4d \mathcal{N} = 3$ SCFTs

Solving $\mathcal{N} = 3$ SCFTs?

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

$$q(z, \bar{z}) = u_I Q' = Q(z, \bar{z}) + \bar{z} \tilde{Q}^*(z, \bar{z}),$$

$$\tilde{q}(z, \bar{z}) = u_I \tilde{Q}' = \tilde{Q}(z, \bar{z}) - \bar{z} Q^*(z, \bar{z})$$

Chiral algebra

Example: free hypermultiplet

Complex scalars in hypermultiplet are in the cohomology

$$Q' = \begin{bmatrix} Q \\ \tilde{Q}^* \end{bmatrix}, \quad \tilde{Q}' = \begin{bmatrix} \tilde{Q} \\ -Q^* \end{bmatrix}$$

$$u_I = (1, \bar{z})$$

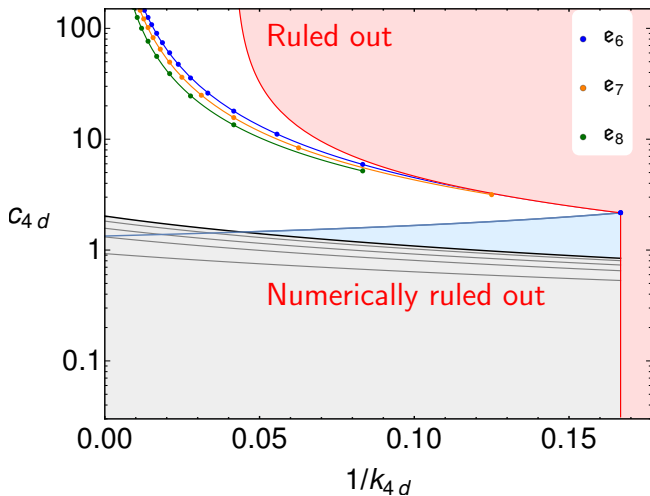
$$q(z, \bar{z}) = u_I Q' = Q(z, \bar{z}) + \bar{z} \tilde{Q}^*(z, \bar{z}),$$

$$\tilde{q}(z, \bar{z}) = u_I \tilde{Q}' = \tilde{Q}(z, \bar{z}) - \bar{z} Q^*(z, \bar{z})$$

$$\rightarrow q(z, z\bar{b})\tilde{q}(0) \sim \bar{z} \tilde{Q}^*(z, \bar{z}) \tilde{Q}(0) \sim \frac{\bar{z}}{z\bar{z}} = \frac{1}{z}$$

$4d \mathcal{N} = 2$ SCFTs with E_6 flavor symmetry

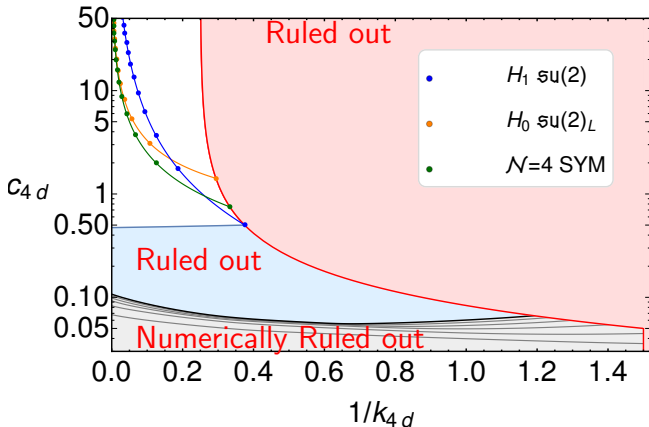
$$\langle TTTT \rangle, \langle J^a J^b J^c J^d \rangle, \langle TTJ^a J^b \rangle$$



What is the space of consistent SCFTs?

4d $\mathcal{N} = 2$ SCFTs with $SU(2)$ flavor symmetry

$\langle TTTT \rangle$, $\langle J^a J^b J^c J^d \rangle$, $\langle TTJ^a J^b \rangle$



Outline

Numerical conformal Bootstrap review

Chiral algebra

$4d \mathcal{N} = 3$ SCFTs

Solving $\mathcal{N} = 3$ SCFTs?

$4d \mathcal{N} = 3$ SCFTs

→ *Non-trivial* interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

$4d \mathcal{N} = 3$ SCFTs

→ *Non-trivial* interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

→ Non-Lagrangian

$4d \mathcal{N} = 3$ SCFTs

→ *Non-trivial* interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

→ Non-Lagrangian

Properties from representation theory [Aharony, Evtikhiev]

$4d \mathcal{N} = 3$ SCFTs

→ *Non-trivial* interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

→ Non-Lagrangian

Properties from representation theory [Aharony, Evtikhiev]

→ $SU(3)_R \times U(1)_r$

$4d \mathcal{N} = 3$ SCFTs

→ *Non-trivial* interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

→ Non-Lagrangian

Properties from representation theory [Aharony, Evtikhiev]

→ $SU(3)_R \times U(1)_r$

→ No flavor symmetry

4d $\mathcal{N} = 3$ SCFTs

→ *Non-trivial* interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

→ Non-Lagrangian

Properties from representation theory [Aharony, Evtikhiev]

→ $SU(3)_R \times U(1)_r$

→ No flavor symmetry

→ $c = a$

4d $\mathcal{N} = 3$ SCFTs

→ *Non-trivial* interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

→ Non-Lagrangian

Properties from representation theory [Aharony, Evtikhiev]

→ $SU(3)_R \times U(1)_r$

→ No flavor symmetry

→ $c = a$

→ No exactly marginal deformations

4d $\mathcal{N} = 3$ SCFTs

→ *Non-trivial* interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

→ Non-Lagrangian

Properties from representation theory [Aharony, Evtikhiev]

→ $SU(3)_R \times U(1)_r$

→ No flavor symmetry

→ $c = a$

→ No exactly marginal deformations

→ Just another SCFT

4d $\mathcal{N} = 3$ SCFTs

→ *Non-trivial* interacting theories

[García-Etxebarria, Regalado] [Aharony, Tachikawa]

→ Non-Lagrangian

Properties from representation theory [Aharony, Evtikhiev]

→ $SU(3)_R \times U(1)_r$

→ No flavor symmetry

→ $c = a$

→ No exactly marginal deformations

→ Just another SCFT

as $\mathcal{N} = 2$: $SU(2)_R \times U(1)_r \times U(1)_F$

Outline

Numerical conformal Bootstrap review

Chiral algebra

$4d \mathcal{N} = 3$ SCFTs

Solving $\mathcal{N} = 3$ SCFTs?

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

→ Set of $\frac{1}{2}$ -BPS operators

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
- Stress tensor supermultiplet
- ▶ Can one write a consistent operator product algebra?

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
- Stress tensor supermultiplet
- ▶ Can one write a consistent operator product algebra?
- ▶ Simplest known theory:

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
- Stress tensor supermultiplet
- ▶ Can one write a consistent operator product algebra?
- ▶ Simplest known theory:
 - $\frac{1}{2}$ -BPS operators of dimension three
 - Stress tensor

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
- Stress tensor supermultiplet
- ▶ Can one write a consistent operator product algebra?
- ▶ Simplest known theory:
 - $\frac{1}{2}$ -BPS operators of dimension three
 - Stress tensor
 - chiral algebra fully fixed and $c_{4d} = \frac{15}{12}$ [Nishinaka, Tachikawa]

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
- Stress tensor supermultiplet
- ▶ Can one write a consistent operator product algebra?
- ▶ Simplest known theory:
 - $\frac{1}{2}$ -BPS operators of dimension three
 - Stress tensor
 - chiral algebra fully fixed and $c_{4d} = \frac{15}{12}$ [Nishinaka, Tachikawa]
 - can compute additional protected OPE coefficients

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

→ Set of $\frac{1}{2}$ -BPS operators

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
- Stress tensor supermultiplet
- ▶ What about higher rank theories?

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
- Stress tensor supermultiplet
- ▶ What about higher rank theories?
 - This would be a closed subalgebra, but c_{4d} is different

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
- Stress tensor supermultiplet
- ▶ What about higher rank theories?
 - This would be a closed subalgebra, but c_{4d} is different
 - Assumed set of generators is incomplete for higher rank

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
- Stress tensor supermultiplet
- ▶ What about higher rank theories?
 - This would be a closed subalgebra, but c_{4d} is different
 - Assumed set of generators is incomplete for higher rank
 - Minimum modification: add a single extra generator

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
 - Stress tensor supermultiplet
 - ▶ What about higher rank theories?
 - This would be a closed subalgebra, but c_{4d} is different
 - Assumed set of generators is incomplete for higher rank
 - Minimum modification: add a single extra generator
 - constructed chiral algebra valid for any c_{4d}
- [ML, Liendo, Meneghelli, Mitev]

Chiral algebras for $\mathcal{N} = 3$ SCFTs

Assume set of generators

- Set of $\frac{1}{2}$ -BPS operators
 - Stress tensor supermultiplet
 - ▶ What about higher rank theories?
 - This would be a closed subalgebra, but c_{4d} is different
 - Assumed set of generators is incomplete for higher rank
 - Minimum modification: add a single extra generator
 - constructed chiral algebra valid for any c_{4d}
- [ML, Liendo, Meneghelli, Mitev]
subalgebra of higher rank theories?

Solving $\mathcal{N} = 3$ SCFTs?

Can we “zoom in” to the $c_{4d} = \frac{15}{12}$ SCFT?

Solving $\mathcal{N} = 3$ SCFTs?

Can we “zoom in” to the $c_{4d} = \frac{15}{12}$ SCFT?

- ▶ $\frac{1}{2}$ -BPS operator of dimension three

Solving $\mathcal{N} = 3$ SCFTs?

Can we “zoom in” to the $c_{4d} = \frac{15}{12}$ SCFT?

- ▶ $\frac{1}{2}$ -BPS operator of dimension three
- ▶ Step 1: Fix protected contributions

Solving $\mathcal{N} = 3$ SCFTs?

Can we “zoom in” to the $c_{4d} = \frac{15}{12}$ SCFT?

- ▶ $\frac{1}{2}$ -BPS operator of dimension three
- ▶ Step 1: Fix protected contributions

$$\lambda_{2d}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2 - \lambda_{\mathcal{O}'_{4d}}^2$$

Solving $\mathcal{N} = 3$ SCFTs?

Can we “zoom in” to the $c_{4d} = \frac{15}{12}$ SCFT?

- ▶ $\frac{1}{2}$ -BPS operator of dimension three
- ▶ Step 1: Fix protected contributions

$$\lambda_{2d}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2 - \lambda_{\mathcal{O}'_{4d}}^2$$

Valid for any $\mathcal{N} = 3$ SCFT with this operator

Solving $\mathcal{N} = 3$ SCFTs?

Can we “zoom in” to the $c_{4d} = \frac{15}{12}$ SCFT?

- ▶ $\frac{1}{2}$ -BPS operator of dimension three
- ▶ Step 1: Fix protected contributions

$$\lambda_{2d}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2 - \lambda_{\mathcal{O}'_{4d}}^2$$

Valid for any $\mathcal{N} = 3$ SCFT with this operator

- ▶ $\mathcal{N} = 4$ “sits in the way”

Solving $\mathcal{N} = 3$ SCFTs?

Can we “zoom in” to the $c_{4d} = \frac{15}{12}$ SCFT?

- ▶ $\frac{1}{2}$ -BPS operator of dimension three
- ▶ Step 1: Fix protected contributions

$$\lambda_{2d}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2 - \lambda_{\mathcal{O}'_{4d}}^2$$

Valid for any $\mathcal{N} = 3$ SCFT with this operator

- ▶ $\mathcal{N} = 4$ “sits in the way”
- ▶ Input chiral algebra data of specific theory

Solving $\mathcal{N} = 3$ SCFTs?

Can we “zoom in” to the $c_{4d} = \frac{15}{12}$ SCFT?

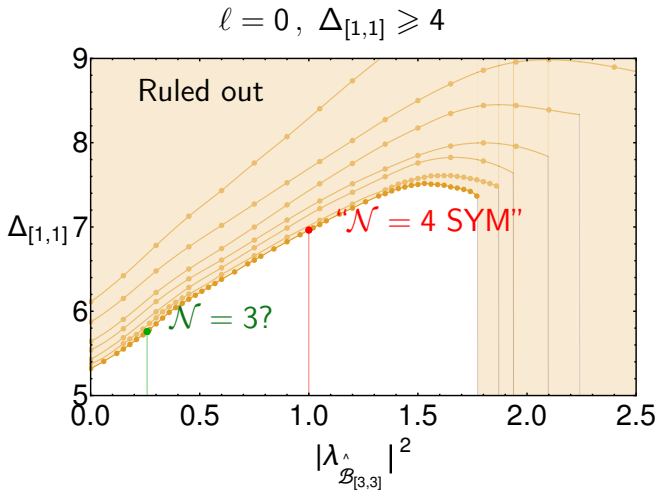
- ▶ $\frac{1}{2}$ -BPS operator of dimension three
- ▶ Step 1: Fix protected contributions

$$\lambda_{2d}^2 \rightsquigarrow \lambda_{\mathcal{O}_{4d}}^2 - \lambda_{\mathcal{O}'_{4d}}^2$$

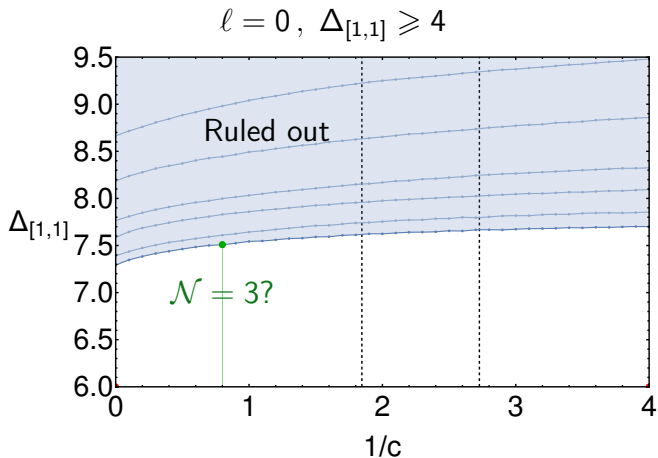
Valid for any $\mathcal{N} = 3$ SCFT with this operator

- ▶ $\mathcal{N} = 4$ “sits in the way”
- ▶ Input chiral algebra data of specific theory
chiral algebra conjectured by [\[Nishinaka, Tachikawa\]](#)

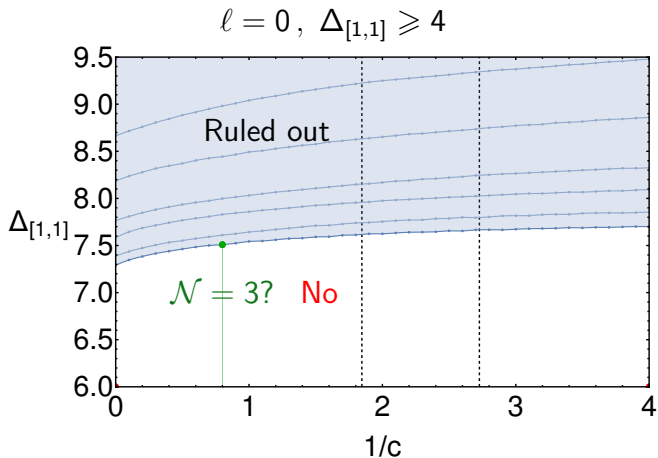
$4d \mathcal{N} = 3$ SCFT with $c = \frac{15}{12}$



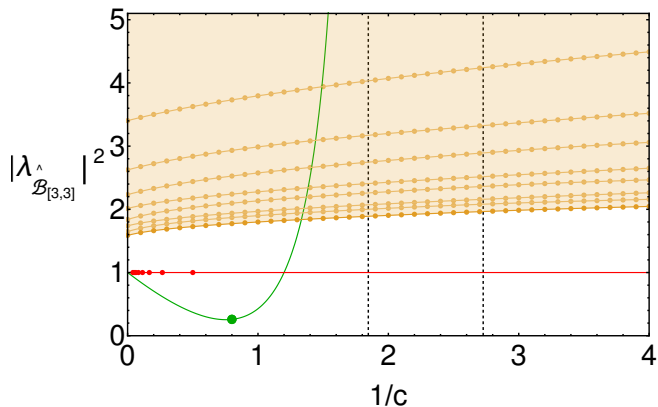
4d $\mathcal{N} = 3$ SCFTs



4d $\mathcal{N} = 3$ SCFTs



4d $\mathcal{N} = 3$ SCFTs



[ML, Liendo, Meneghelli, Mitev]