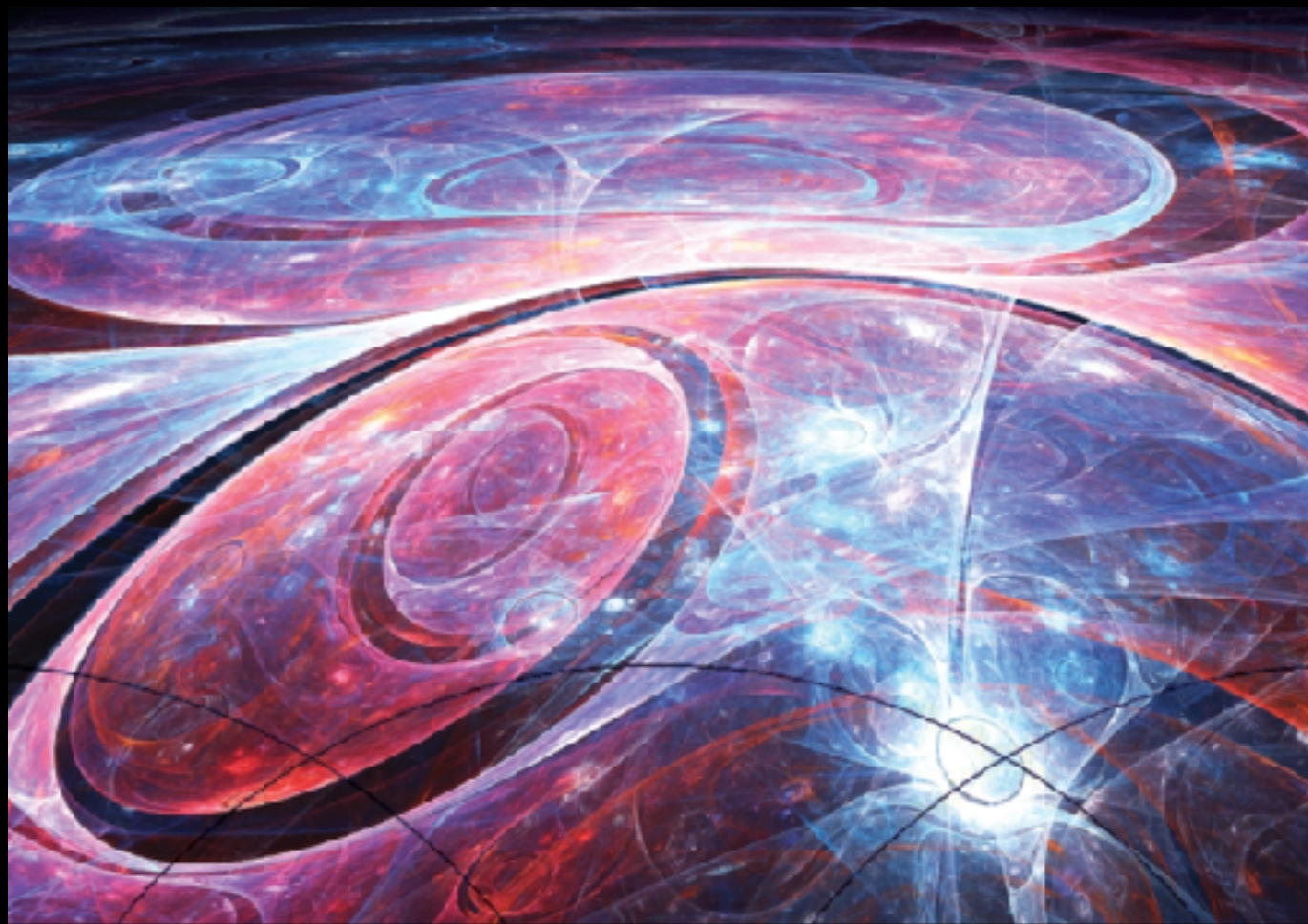


BPS-states and automorphic representations of exceptional groups

Daniel Persson
Chalmers University of Technology



String Math
Hamburg
July 26, 2017

Outline

1. Motivation from string theory

2. Automorphic forms and representation theory

3. Small representations and BPS-couplings

4. Outlook



Talk based on our papers/book:

[1511.04265] w/ Fleig, Gustafsson, Kleinschmidt

[1412.5625] w/ Gustafsson, Kleinschmidt

[1312.3643] w/ Fleig, Kleinschmidt

to appear on Friday this week

[1707.XXXX] w/ Ahlén, Gustafsson, Kleinschmidt, Liu

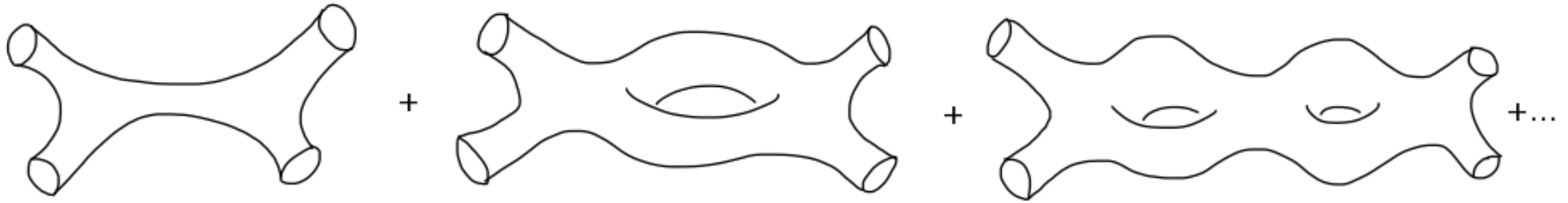
and in progress

[17YY.XXXX] w/ Gourevitch, Gustafsson, Kleinschmidt, Sahi

I. Motivation from string theory

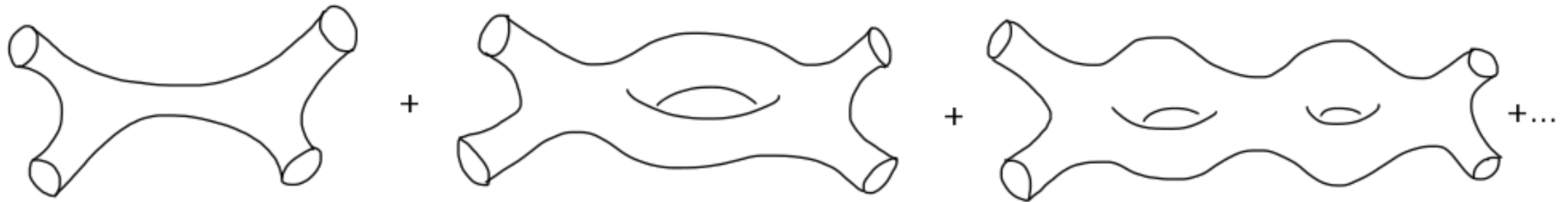
String amplitudes

Understand the structure of **string interactions**



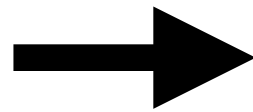
String amplitudes

Understand the structure of **string interactions**



Strongly constrained by **symmetries!**

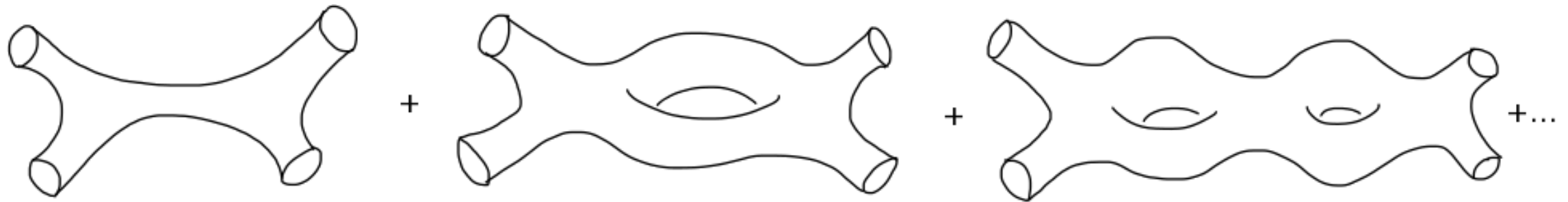
- supersymmetry
- U-duality



amplitudes have intricate
arithmetic structure $G(\mathbb{Z})$

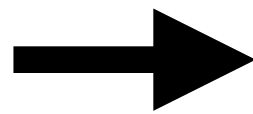
String amplitudes

Understand the structure of **string interactions**



Strongly constrained by **symmetries**!

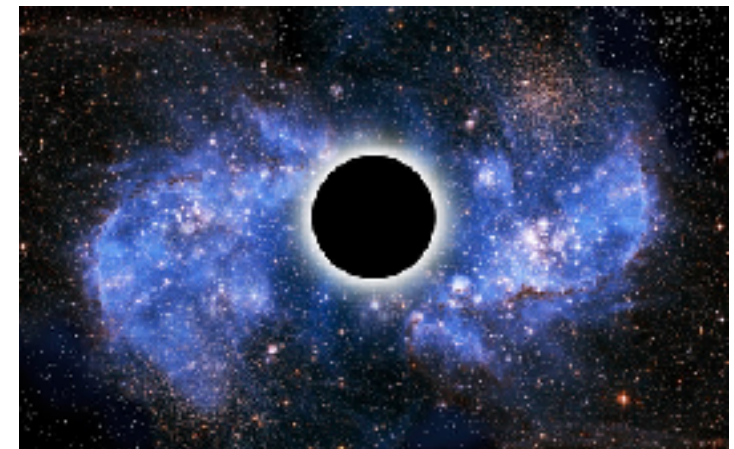
- supersymmetry
- U-duality



amplitudes have intricate **arithmetic structure** $G(\mathbb{Z})$

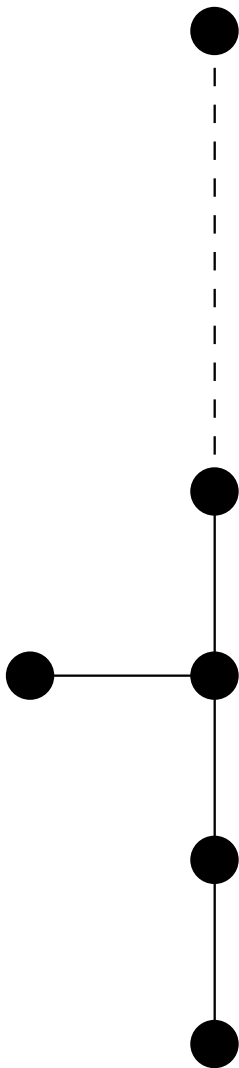
Symmetry constrains interactions, leads to insights about:

- ultraviolet properties of gravity
- non-perturbative effects (black holes, instantons)
- novel mathematical predictions from physics



Toroidal compactifications yield the famous chain of **U-duality groups**

[Cremmer, Julia][Hull, Townsend]



D	G	K	$G(\mathbb{Z})$
10	$SL(2, \mathbb{R})$	$SO(2)$	$SL(2, \mathbb{Z})$
9	$SL(2, \mathbb{R}) \times \mathbb{R}^+$	$SO(2)$	$SL(2, \mathbb{Z})$
8	$SL(3, \mathbb{R}) \times SL(2, \mathbb{R})$	$SO(3) \times SO(2)$	$SL(3, \mathbb{Z}) \times SL(2, \mathbb{Z})$
7	$SL(5, \mathbb{R})$	$SO(5)$	$SL(5, \mathbb{Z})$
6	$Spin(5, 5, \mathbb{R})$	$(Spin(5) \times Spin(5))/\mathbb{Z}_2$	$Spin(5, 5, \mathbb{Z})$
5	$E_6(\mathbb{R})$	$USp(8)/\mathbb{Z}_2$	$E_6(\mathbb{Z})$
4	$E_7(\mathbb{R})$	$SU(8)/\mathbb{Z}_2$	$E_7(\mathbb{Z})$
3	$E_8(\mathbb{R})$	$Spin(16)/\mathbb{Z}_2$	$E_8(\mathbb{Z})$

Physical couplings are given by **automorphic forms** on


$$G(\mathbb{Z}) \backslash G(\mathbb{R}) / K$$

Green, Gutperle, Sethi, Vanhove, Kiritsis, Pioline, Obers, Kazhdan, Waldron, Basu, Russo, Cederwall, Bao, Nilsson, D.P., Lambert, West, Gubay, Miller, Bossard, Fleig, Kleinschmidt, Gustafsson, Cosnier-Horeau...

Higher-derivative action in type II string theory on tori

$$\int d^{10-n}x \sqrt{G} \left[(\alpha')^3 f_0(g) \mathcal{R}^4 + (\alpha')^5 f_4(g) \partial^4 \mathcal{R}^4 + \cdots \right]$$

Higher-derivative action in type II string theory on tori

$$\int d^{10-n}x \sqrt{G} \left[(\alpha')^3 f_0(g) \mathcal{R}^4 + (\alpha')^5 f_4(g) \partial^4 \mathcal{R}^4 + \cdots \right]$$


contraction of four Riemann tensors

Higher-derivative action in type II string theory on tori

$$\int d^{10-n}x \sqrt{G} \left[(\alpha')^3 f_0(g) \mathcal{R}^4 + (\alpha')^5 f_4(g) \partial^4 \mathcal{R}^4 + \cdots \right]$$

- $f_0(g), f_4(g)$ are functions of $g \in E_{n+1}(\mathbb{R})/K$
- must be **invariant** under U-duality $E_{n+1}(\mathbb{Z})$
- supersymmetry requires that they are **Laplacian eigenfunctions**
- well-defined **weak-coupling expansions** as $g_s \rightarrow 0$

Higher-derivative action in type II string theory on tori

$$\int d^{10-n}x \sqrt{G} \left[(\alpha')^3 f_0(g) \mathcal{R}^4 + (\alpha')^5 f_4(g) \partial^4 \mathcal{R}^4 + \cdots \right]$$

- $f_0(g), f_4(g)$ are functions of $g \in E_{n+1}(\mathbb{R})/K$
- must be **invariant** under U-duality $E_{n+1}(\mathbb{Z})$
- supersymmetry requires that they are **Laplacian eigenfunctions**
- well-defined **weak-coupling expansions** as $g_s \rightarrow 0$

defining properties
of an
**automorphic
form!**

2. Automorphic forms and representation theory

Data:

- ▶ $G(\mathbb{R})$ real semi-simple Lie group (e.g. $SL(n, \mathbb{R})$)
- ▶ $G(\mathbb{Z}) \subset G$ arithmetic subgroup (e.g. $SL(n, \mathbb{Z})$)

Data:

- ▶ $G(\mathbb{R})$ real semi-simple Lie group (e.g. $SL(n, \mathbb{R})$)
- ▶ $G(\mathbb{Z}) \subset G$ arithmetic subgroup (e.g. $SL(n, \mathbb{Z})$)

Definition:

An **automorphic form** is a smooth function $\varphi : G \longrightarrow \mathbb{C}$ satisfying

1. Automorphy: $\forall \gamma \in G(\mathbb{Z}), \varphi(\gamma g) = \varphi(g)$
2. φ is an eigenfunction of the ring of inv. diff. operators on G
3. φ has well-behaved growth conditions

Data:

- ▶ $G(\mathbb{R})$ real semi-simple Lie group (e.g. $SL(n, \mathbb{R})$)
- ▶ $G(\mathbb{Z}) \subset G$ arithmetic subgroup (e.g. $SL(n, \mathbb{Z})$)

Example: Non-holomorphic Eisenstein series on $G(\mathbb{R}) = SL(2, \mathbb{R})$

$$E(s, \tau) = \sum_{\substack{(m,n) \in \mathbb{Z}^2 \\ (m,n) \neq (0,0)}} \frac{y^s}{|m\tau + n|^{2s}} \quad s \in \mathbb{C}$$

$$\tau = x + iy \in \mathbb{H} \cong SL(2, \mathbb{R})/U(1)$$

Automorphic representations

$$\mathcal{A}(G(\mathbb{Z}) \backslash G(\mathbb{R})) = \{\text{space of automorphic forms on } G(\mathbb{R})\}$$

Automorphic representations

$$\mathcal{A}(G(\mathbb{Z}) \backslash G(\mathbb{R})) = \{\text{space of automorphic forms on } G(\mathbb{R})\}$$

The group G acts on this space via the **right-regular representation**:

$$(\rho(h)\varphi)(g) = \varphi(gh)$$

for $\varphi \in \mathcal{A}$ and $h, g \in G$

Automorphic representations

$\mathcal{A}(G(\mathbb{Z}) \backslash G(\mathbb{R})) = \{\text{space of automorphic forms on } G(\mathbb{R})\}$

The group G acts on this space via the **right-regular representation**:

$$(\rho(h)\varphi)(g) = \varphi(gh)$$

for $\varphi \in \mathcal{A}$ and $h, g \in G$

Definition: An **automorphic representation** π of G is an irreducible constituent in the decomposition of \mathcal{A} under the right-regular action.

Automorphic representations

$$\mathcal{A}(G(\mathbb{Z}) \backslash G(\mathbb{R})) = \mathcal{A}_{discrete} \oplus \mathcal{A}_{continuous}$$

➔ $\mathcal{A}_{discrete}$: generated by cusp forms
(and residues of Eisenstein series)

$$\int_{U(\mathbb{Z}) \backslash U(\mathbb{R})} \varphi(ug) du = 0 \quad \begin{array}{l} \text{all unipotents} \\ U \subset G \end{array}$$

➔ $\mathcal{A}_{continuous}$: generated by Eisenstein series

Example: $G = SL(2, \mathbb{R})$

$$\int d^{10}x \sqrt{G} f_0(\tau) \mathcal{R}^4$$

$$y = g_s^{-1}$$

perturbative terms

non-perturbative terms

$$f_0(\tau) = \underbrace{2\zeta(3)y^{3/2}}_{\text{tree-level}} + \underbrace{4\zeta(2)y^{-1/2}}_{\text{one-loop}} + 2\pi \sum_{m \neq 0} \sqrt{|m|} \sigma_{-2}(m) e^{-S_{\text{inst}}(z)} [1 + \mathcal{O}(y^{-1})]$$

tree-level

one-loop



amplitudes in the presence of instantons

**unique
solution!**

$$f_0(\tau) = \sum_{(m,n) \neq (0,0)} \frac{y^{3/2}}{|m + n\tau|^3}$$

[Green, Gutperle]
[Green, Sethi]
[Pioline]

Example: $G = SL(2, \mathbb{R})$

$$\int d^{10}x \sqrt{G} f_0(\tau) \mathcal{R}^4$$

$$y = g_s^{-1}$$

perturbative terms

non-perturbative terms

$$f_0(\tau) = \underbrace{2\zeta(3)y^{3/2}}_{\text{tree-level}} + \underbrace{4\zeta(2)y^{-1/2}}_{\text{one-loop}} + 2\pi \sum_{m \neq 0} \sqrt{|m|} \sigma_{-2}(m) e^{-S_{\text{inst}}(z)} [1 + \mathcal{O}(y^{-1})]$$

tree-level

one-loop



amplitudes in the presence of instantons

instanton action

$$S_{\text{inst}}(z) := 2\pi |m| y - 2\pi i m x$$

Example: $G = SL(2, \mathbb{R})$

$$\int d^{10}x \sqrt{G} f_0(\tau) \mathcal{R}^4$$

$$y = g_s^{-1}$$

perturbative terms

non-perturbative terms

$$f_0(\tau) = \underbrace{2\zeta(3)y^{3/2}}_{\text{tree-level}} + \underbrace{4\zeta(2)y^{-1/2}}_{\text{one-loop}} + 2\pi \sum_{m \neq 0} \sqrt{|m|} \sigma_{-2}(m) e^{-S_{\text{inst}}(z)} [1 + \mathcal{O}(y^{-1})]$$

tree-level

one-loop



amplitudes in the presence of instantons

instanton action

instanton measure

$$S_{\text{inst}}(z) := 2\pi |m| y - 2\pi i m x$$

$$\sigma_{-2}(m) = \sum_{d|m} d^{-2}$$

Adelic framework

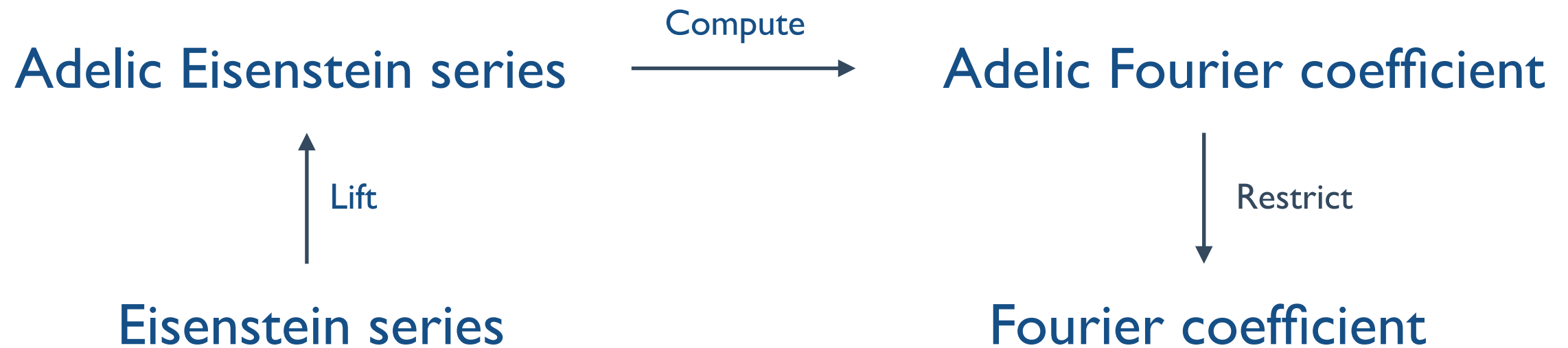
An efficient, but abstract, way to approach the subject of automorphic forms is by the introduction of adeles, rather ungainly objects that nevertheless, once familiar, spare much unnecessary thought and many useless calculations.

— Robert P. Langlands

Adelic framework

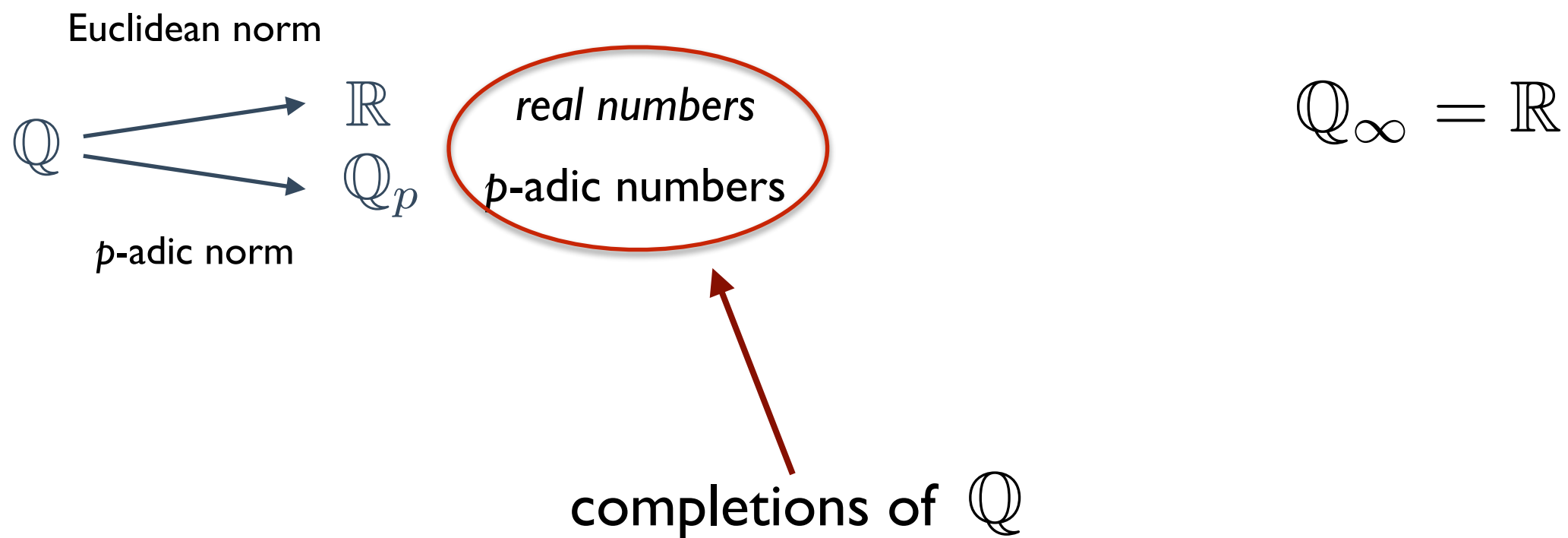
An efficient, but abstract, way to approach the subject of automorphic forms is by the introduction of adeles, rather ungainly objects that nevertheless, once familiar, spare much unnecessary thought and many useless calculations.

— Robert P. Langlands



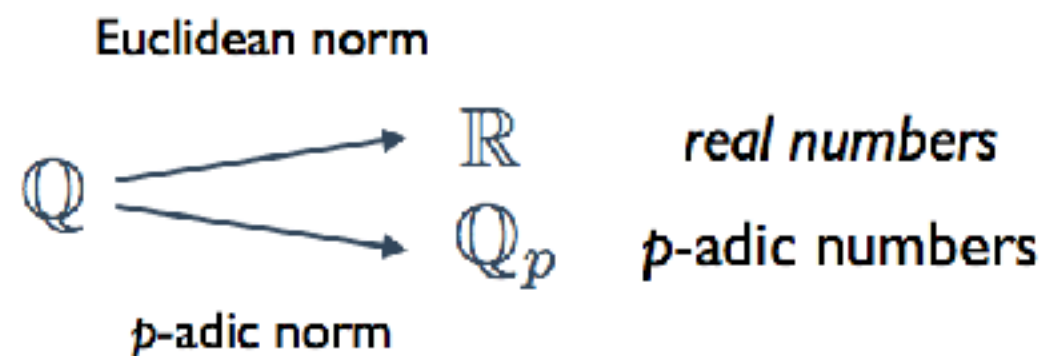
Adelic framework

For each prime number p



Adelic framework

For each prime number p



$$\mathbb{Q}_\infty = \mathbb{R}$$

The adeles are then defined as

$$\mathbb{A} = \mathbb{R} \times \prod'_{p \text{ prime} < \infty} \mathbb{Q}_p$$

global **local**

$$x = (x_\infty; x_2, x_3, x_5, \dots) \in \mathbb{A}$$

$$\mathbb{Q} \hookrightarrow \mathbb{A}$$

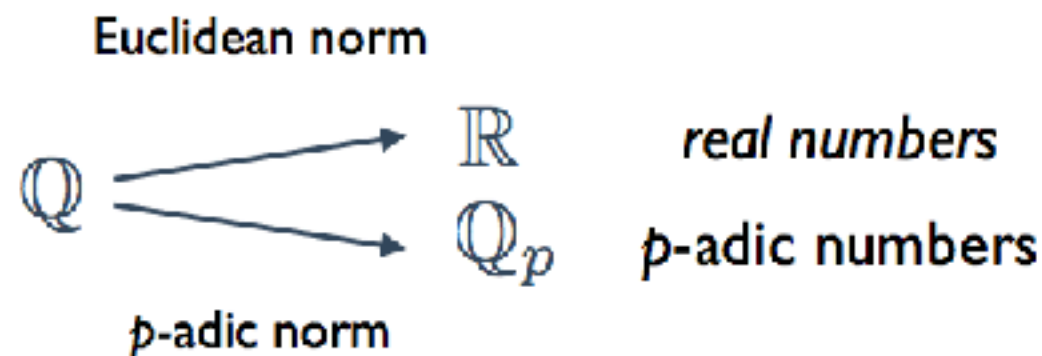
$$\mathbb{Q} \subset \mathbb{A} \quad \text{discrete embedding}$$

$$q \mapsto (q; q, q, q, \dots)$$

analogous to: $\mathbb{Z} \subset \mathbb{R}$

Adelic framework

For each prime number p



$$\mathbb{Q}_\infty = \mathbb{R}$$

The adeles are then defined as

$$\mathbb{A} = \mathbb{R} \times \prod'_{p \text{ prime} < \infty} \mathbb{Q}_p$$

global **local**

$x = (x_\infty; x_2, x_3, x_5, \dots) \in \mathbb{A}$

$$\mathbb{Q} \hookrightarrow \mathbb{A}$$

$$\mathbb{Q} \subset \mathbb{A}$$

much easier to work with
since \mathbb{Q} is a **field**.

$$q \mapsto (q; q, q, q, \dots)$$

analogous to: $\mathbb{Z} \subset \mathbb{R}$

Adelic framework

(completed) **Riemann zeta function:** $\xi(s) = \xi(1-s)$

$$\xi(s) = \pi^{-s/2} \Gamma(s/2) \zeta(s) = \pi^{-s/2} \Gamma(s/2) \prod_{p \text{ prime} < \infty} \frac{1}{1-p^{-s}}$$

$$= \int_{\mathbb{R}} e^{-\pi x^2} |x|^s dx \prod_{p \text{ prime} < \infty} \int_{\mathbb{Q}_p} \gamma_p(x) |x|_p^s dx$$

$$= \int_{\mathbb{A}} \gamma_{\mathbb{A}}(x) |x|_{\mathbb{A}}^s dx$$

In his famous thesis, Tate gave elegant new proofs of the **functional equation and analytic continuation** of $\xi(s)$ using these techniques

Adelic framework

arithmetic groups

$$G(\mathbb{Z}) \subset G(\mathbb{R}) \longrightarrow G(\mathbb{Q}) \subset G(\mathbb{A})$$

space of
automorphic forms

$$\mathcal{A}(G(\mathbb{Z}) \backslash G(\mathbb{R})) \longrightarrow \mathcal{A}(G(\mathbb{Q}) \backslash G(\mathbb{A}))$$

Adelic framework

arithmetic groups

$$G(\mathbb{Z}) \subset G(\mathbb{R}) \longrightarrow G(\mathbb{Q}) \subset G(\mathbb{A})$$

space of automorphic forms

$$\mathcal{A}(G(\mathbb{Z}) \backslash G(\mathbb{R})) \longrightarrow \mathcal{A}(G(\mathbb{Q}) \backslash G(\mathbb{A}))$$

$$\cup$$

$$\cup$$

Eisenstein series

$$\sum_{\gamma \in B(\mathbb{Z}) \backslash G(\mathbb{Z})} e^{\langle \lambda + \rho | H(\gamma g) \rangle} \longrightarrow \sum_{\gamma \in B(\mathbb{Q}) \backslash G(\mathbb{Q})} e^{\langle \lambda + \rho | H(\gamma g) \rangle}$$

$$\lambda \in \mathfrak{h}^* \otimes \mathbb{C}$$

$$H : G \rightarrow \mathfrak{h}$$

3. Small representations and BPS-couplings

Minimal automorphic representations

Definition: *An automorphic representation*

$$\pi = \bigotimes_{p \leq \infty} \pi_p$$

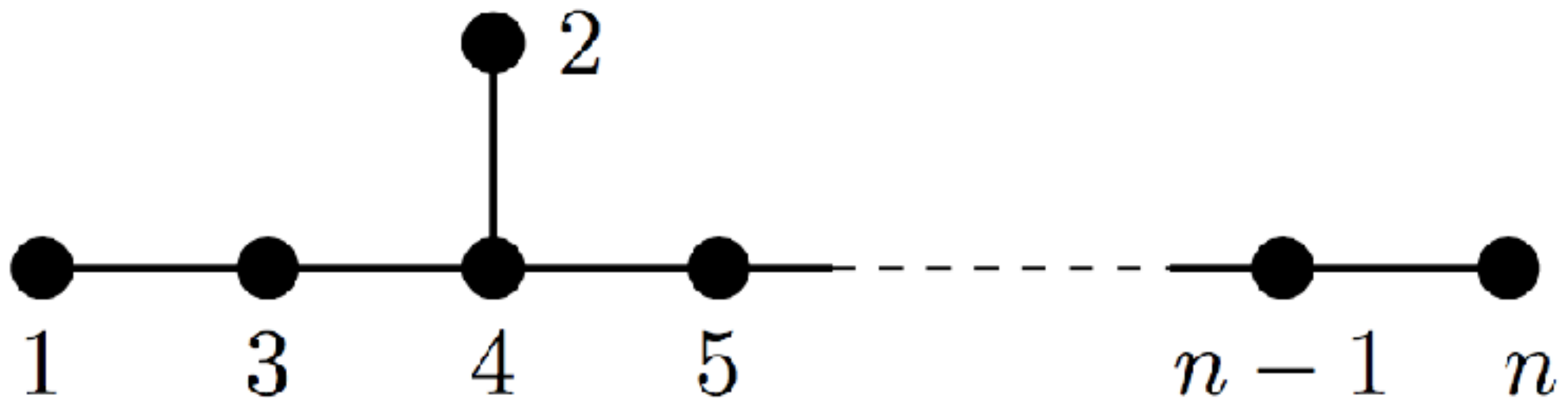
is minimal if each factor π_p has smallest non-trivial Gelfand-Kirillov dimension.

[Joseph][Brylinski, Kostant][Ginzburg, Rallis, Soudry][Kazhdan, Savin]....

Automorphic forms $\varphi \in \pi_{min}$ are characterised by having
very few non-vanishing Fourier coefficients.

[Ginzburg, Rallis, Soudry]

Exceptional groups



Functional dimension of minimal representations:

$$\text{GKdim } \pi_{\min} = \begin{cases} 11, & E_6 \\ 17, & E_7 \\ 29, & E_8 \end{cases}$$

Automorphic realization

Consider the Borel-Eisenstein series on $G(\mathbb{A})$

$$E(\lambda, g) = \sum_{\gamma \in B(\mathbb{Q}) \backslash G(\mathbb{Q})} e^{\langle \lambda + \rho | H(\gamma g) \rangle}$$

Now fix the weight to

$$\lambda = 2s\Lambda_1 - \rho$$

where Λ_1 is the fundamental weight associated to node 1.

Theorem [Ginzburg,Rallis,Soudry][Green,Miller,Vanhove]

For $G = E_6, E_7, E_8$ the Eisenstein series $E(2s\Lambda - \rho, g)$ evaluated at $s = 3/2$ is attached to the representation π_{min}

This theorem yields an explicit automorphic realisation of the minimal representation.

Our aim is to use this to calculate Fourier coefficients associated with maximal parabolic subgroups.

BPS-couplings

$$g \in E_n(\mathbb{R})$$

$$\int d^{11-n} x \sqrt{G} f_0(g) \mathcal{R}^4 \quad f_0(g) = E(3/2, g) \quad s = 3/2$$

$$\int d^{11-n} \sqrt{G} f_4(g) \partial^4 \mathcal{R}^4 \quad f_4(g) = E(5/2, g) \quad s = 5/2$$

These partition functions are Eisenstein series attached to **small automorphic representations** of G .

[Green, Miller, Vanhove][Pioline]

minimal automorphic
representation

$$\pi_{min}$$

1/2 - BPS

next-to-minimal automorphic
representation

$$\pi_{ntm}$$

1/4 - BPS

Perturbative limit - choices of unipotent subgroups

→ **Decompactification limit**



- perturbative parameter: radius of decompactified circle
- non-perturbative effects: KK-instantons, BPS-instantons

→ **String perturbation limit**



- perturbative parameter: string coupling
- non-perturbative effects: D-instantons, NS5-instantons

→ **M-theory limit**



- perturbative parameter: volume of M-theory torus
- non-perturbative effects: M2- & M5-instantons

general Fourier coefficients

- ▶ $P = LU$ **standard parabolic** of G
- ▶ **unitary character** $\psi_U : U(\mathbb{Q}) \backslash U(\mathbb{A}) \rightarrow U(1)$
- ▶ We then have the U -**Fourier coefficient**:

$$F_{\psi_U}(f_\chi, g) = \int_{U(\mathbb{Q}) \backslash U(\mathbb{A})} E(f_\chi, ug) \overline{\psi_U(u)} du$$

very little is known in general in this case...

Theorem [Miller-Sahi]: Let G be a split group of type E_6 or E_7
Then any Fourier coefficient of $\varphi \in \pi_{min}$ of G is completely
determined by the maximally degenerate Whittaker coefficient

$$W_{\psi_\alpha}(\varphi, g) = \int_{N(\mathbb{Q}) \backslash N(\mathbb{A})} \varphi(n g) \overline{\psi_\alpha(n)} dn$$

Theorem [Miller-Sahi]: Let G be a split group of type E_6 or E_7
 Then any Fourier coefficient of $\varphi \in \pi_{min}$ of G is completely
 determined by the maximally degenerate Whittaker coefficient

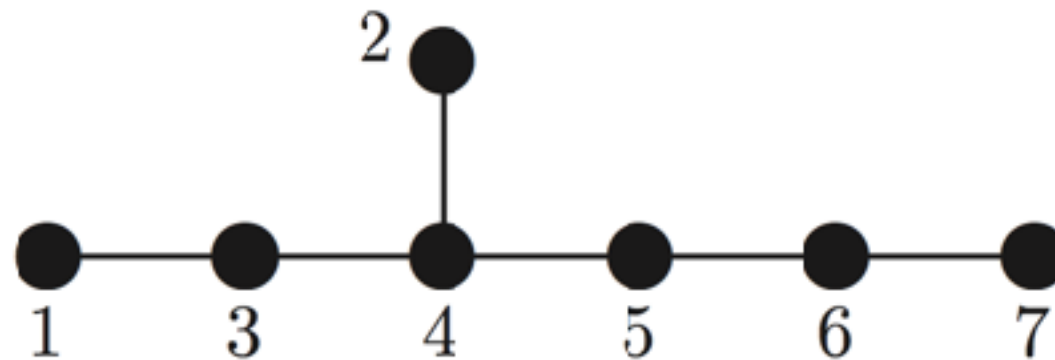
$$W_{\psi_\alpha}(\varphi, g) = \int_{N(\mathbb{Q}) \backslash N(\mathbb{A})} \varphi(ng) \overline{\psi_\alpha(n)} dn$$

Can one use this to calculate

$$F_{\psi_U}(\varphi, g) = \int_{U(\mathbb{Q}) \backslash U(\mathbb{A})} E(\varphi, ug) \overline{\psi_U(u)} du$$

in terms of W_{ψ_α} ?

Example: $G = E_7$



Consider the **3-grading** of the Lie algebra

$$\mathfrak{e}_7 = \mathfrak{g}_{-1} \oplus \mathfrak{g}_0 \oplus \mathfrak{g}_1 = \mathbf{27} \oplus (\mathfrak{e}_6 \oplus \mathbf{1}) \oplus \mathbf{27}$$

The space $\mathfrak{g}_0 \oplus \mathfrak{g}_1$ is the Lie algebra of a maximal parabolic $P = LU$ with 27-dim unipotent U and Levi $L = E_6 \times GL(1)$

The degenerate Whittaker vector associated with α_7 is given by:
[Fleig, Kleinschmidt, D.P.]

$$W_{\psi_k}(3/2, a) = |k|^{3/2} \sigma_{-3}(k) K_{3/2}(2\pi |k| a)$$

where $a \in A \subset E_7$ and

$$\sigma_s(k) = \sum_{d|k} d^s$$

The degenerate Whittaker vector associated with α_7 is given by:
[Fleig, Kleinschmidt, D.P.]

$$W_{\psi_k}(3/2, a) = |k|^{3/2} \sigma_{-3}(k) K_{3/2}(2\pi |k| a)$$

where $a \in A \subset E_7$ and

$$\sigma_s(k) = \sum_{d|k} d^s$$

We now want to relate this to the U - Fourier coefficient

$$F_{\psi_U}(3/2, g) = \int_{U(\mathbb{Q}) \backslash U(\mathbb{A})} E(3/2, ug) \overline{\psi_U(u)} du$$

This captures **instantons in the decompactification limit** of II/T^6 !

Claim: [\[Pioline\]](#) [\[Gustafsson, Kleinschmidt, D.P.\]](#) [\[Bossard, Vershinin\]](#)

$$F_{\psi_U}(3/2; h, r) = |k|^{3/2} \sigma_{-3}(k) K_{3/2}(2\pi r |k| \times ||h^{-1} \vec{x}||)$$

where $h \in E_6$, $r \in GL(1)$ and $\vec{x} \in \mathbb{Z}^{27}$

Claim: [\[Pioline\]](#) [\[Gustafsson, Kleinschmidt, D.P.\]](#) [\[Bossard, Verschinin\]](#)

$$F_{\psi_U}(3/2; h, r) = |k|^{3/2} \sigma_{-3}(k) K_{3/2}(2\pi r |k| \times ||h^{-1} \vec{x}||)$$

where $h \in E_6$, $r \in GL(1)$ and $\vec{x} \in \mathbb{Z}^{27}$

Proof: *To appear by* [\[Gourevitch, Gustafsson, Kleinschmidt, D.P., Sahi\]](#)

This gives the **complete abelian Fourier expansion** of the minimal representation

Physically the vector \vec{x} corresponds to the **instanton charges** of the 27 vector fields in D=5.

Next-to-minimal representations

Relevant for $\partial^4 \mathcal{R}^4$ -couplings.

Theorem [Green, Miller, Vanhove]: Let $G = E_6, E_7, E_8$
The Eisenstein series

$$E(s, g) = \sum_{\gamma \in B(\mathbb{Q}) \backslash G(\mathbb{Q})} e^{\langle 2s\Lambda_1 | H(\gamma g) \rangle}$$

evaluated at $s = 5/2$ is a spherical vector in π_{ntm} .

Conjecture [\[Gustafsson, Kleinschmidt, D.P.\]](#):

Let G be a semisimple, simply laced Lie group.

Then all Fourier coefficients of $\varphi \in \pi_{ntm}$ are completely determined by degenerate Whittaker vectors of the form

$$W_{\psi_\alpha}(\varphi, g) = \int_{N(\mathbb{Q}) \backslash N(\mathbb{A})} \varphi(n g) \overline{\psi_\alpha(n)} dn$$

$$W_{\psi_{\alpha, \beta}}(\varphi, g) = \int_{N(\mathbb{Q}) \backslash N(\mathbb{A})} \varphi(n g) \overline{\psi_{\alpha, \beta}(n)} dn$$

where (α, β) are commuting simple roots.

Proof. For $SL(n)$ to appear by [\[Ahlén, Gustafsson, Kleinschmidt, Liu, D.P.\]](#)

For exceptionals, in progress by [\[Gustafsson, Gourevitch, Kleinschmidt, D.P., Sahi\]](#)

Conjecture [\[Gustafsson, Kleinschmidt, D.P.\]](#):

Let G be a semisimple, simply laced Lie group.

Then all Fourier coefficients of $\varphi \in \pi_{ntm}$ are completely determined by degenerate Whittaker vectors of the form

$$W_{\psi_\alpha}(\varphi, g) = \int_{N(\mathbb{Q}) \backslash N(\mathbb{A})} \varphi(ng) \overline{\psi_\alpha(n)} dn$$

$$W_{\psi_{\alpha,\beta}}(\varphi, g) = \int_{N(\mathbb{Q}) \backslash N(\mathbb{A})} \varphi(ng) \overline{\psi_{\alpha,\beta}(n)} dn$$

where (α, β) are commuting simple roots.

This will allow us to **extract instanton effects** from $\partial^4 \mathcal{R}^4$ couplings!

See also [\[Bossard, Pioline\]](#)[\[Bossard, Cosnier-Horeau, Pioline\]](#)

4. Outlook

So what happens at the next order?

1/8-BPS and non-BPS couplings seem to require more general automorphic objects.

[\[Green, Miller, Vanhove\]](#)[\[Pioline\]](#)[\[Bossard, Verschinin\]](#)[\[Bossard, Kleinschmidt\]](#)

This is uncharted
mathematical territory.



So what happens at the next order?

1/8-BPS and non-BPS couplings seem to require more general automorphic objects.

[Green, Miller, Vanhove][Pioline][Bossard, Verschinin][Bossard, Kleinschmidt]

This is uncharted mathematical territory.



Below $D=3$ we also enter the realm of **Kac-Moody groups**!

Connections with **double affine Hecke algebras** (DAHAs)?